

Advecting Procedural Textures for 2D Flow Animation

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Abstract

This paper proposes the use of specially generated 3D procedural textures for visualizing steady state 2D flow fields. We use the flow field to advect and animate the texture over time. However, using standard texture advection techniques and arbitrary textures will introduce some undesirable effects such as: (a) expanding texture from a critical source point, (b) streaking pattern from the boundary of the flow field, (c) crowding of advected textures near an attracting spiral or sink, and (d) absent or lack of textures in some regions of the flow. This paper proposes a number of strategies to solve these problems. We demonstrate how the technique works using both synthetic data and computational fluid dynamics data.

Key Words and Phrases: Procedural textures, critical points, texture advection, animation, vector field.

1 INTRODUCTION

It is hard to avoid staring at the dynamic motions of characters, objects, or almost anything else that moves in an animated scene. Our attention span seems to last longer when watching a video than watching a set of static presentation slides. Recent advances in computer graphics and imaging research have introduced many animation techniques that are based on flow advection [5, 7, 8]. These techniques allow effective visualization of the flow trajectory by observing the movements and distortions of the textures. However, most of them assume simple flow behaviors. In a computational fluid dynamics (CFD) simulation, it is common to produce data sets with very complex flow patterns. Some of these include: saddles, repelling nodes, repelling spirals, attracting nodes, attracting spirals, and centers (circular flows). Using conventional texture advection techniques to animate these types flow may create undesirable artifacts.

A simple method for animating a 2D flow using texture advection can be described as follows: Consider an input

image and a velocity field defined over the image. Let $n + 1$ be the total number of time steps in the animation, then for each time step t_i , where $i = 0, \dots, n$, advect every pixel backward from time t_i to t_0 and save the morphed image M_i . To animate the flow, simply playback the saved images M_i for all time steps. This texture advection method will work for most flow fields. However, when there are repelling nodes or repelling spirals in the flow, then undesirable artifacts can be seen. In flow topology, a repelling node is known as a critical point that generates a star-like outward flow pattern while a repelling spiral is another type of critical point that generates an outward swirl flow pattern. For both types of critical points, the flow would appear to be repelling away from the critical point during flow animation. Hence, they are also known as *source critical points*.

Once the texture at a source critical point is advected, what texture should be used for the next advection at that critical point? Since the flow is being repelled away from the critical point, there is no new texture available! With a texture advection method such as the one just described earlier, one would see an expansion of the original texture color from the source critical point; thus, creating an undesirable artifacts. This artifact can also occur at the image border where there are incoming flows. Analogously, the problem is that once the original texture at the border has been advected forward, what texture should be used next at the border? Here the artifact is caused by an expansion of constant texture from the border. Figure 1 shows a time sequence of animating 2D flows using the texture advection method described above. The left image shows the input texture image. As we step forward in time, we can see the artifacts generated by the source critical point located near the lower-right of the image. The constant bands of texture from the upper-left border of the image are caused by the inflow from the border.

One may consider using periodic textures to feed in textures from the other side as soon as a backward integrated streamline hits a boundary. However, the underlying cause of these artifacts can be attributed to the absence of flow information to guide the texture selection. For the case of

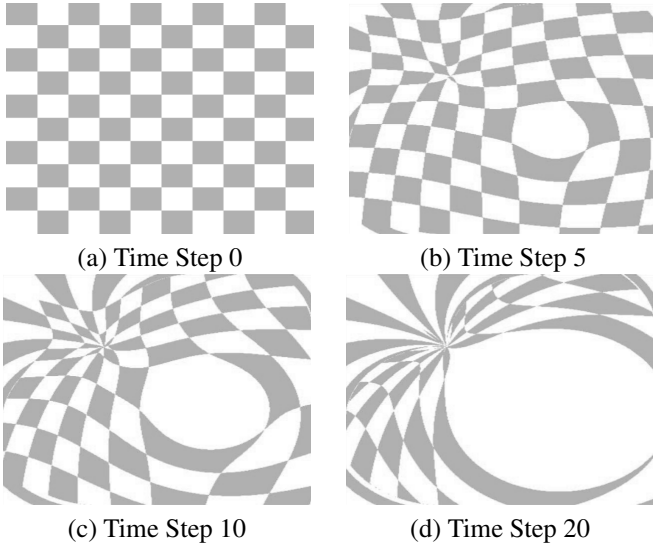


Figure 1. A time sequence of animating 2D flows using texture advection. Artifacts of constant texture due to the source critical point located on the lower right portion of the image and the constant texture bands from the upper left border due to inflow.

inflow from the border, we do not have flow data to determine the next position of the streamline when it is integrated backwards. Likewise, for a source critical point, the backward integrated streamline always arrive at the same point. Hence, the artifact of expansion from a single texel. In our earlier work, we proposed an alternative method to texture advection, based on streamline cycling, to alleviate these artifacts [3]. In this paper, we introduce a technique for eliminating these artifacts using the more traditional texture advection method in conjunction with procedurally generated textures that appear more continuous. We also propose two texture patterns that are ideal for animating 2D flow fields. Finally, we apply our techniques to two real world flow data.

2 FLOWS FROM BOUNDARY

Assume that a steady state vector function $\vec{V}(p)$ is defined for all position p in the image grid G and $t \in [t_0, t_n]$, where $n + 1$ is the total number of time steps in the animation. For texture advection, each pixel in the image is considered as a massless particle. To create an animation, we integrate backwards to determine where the particle at p might have come from. We then sequentially transfer the textures over the path of the particle to p .

The path of particle at p can be computed by numerically

integrating the vector function $\vec{V}(p)$. Since we are dealing with steady flow fields, we can treat the animation time at t_i as our integration time at t_k . That is, $t_k = t_i$. Let t_k be the current time step, p_0 be the initial particle position for each pixel, and initially, $k = 0$, then the particle can be advected backwards from t_k to t_0 using a second-order Runge-Kutta integration with:

$$\begin{aligned}
 K &= -h\vec{V}(p_k) \\
 p_{k+1} &= p_k - h\vec{V}\left(p_k + \frac{K}{2}\right) \\
 t_{k+1} &= t_k - h \quad \text{and} \quad k = k + 1, \quad (1)
 \end{aligned}$$

where h is the integration step size and $0 < h < 1$. The texture at p_0 is then replaced by that at p_k to generate an animation frame for time t_k . We referred to the texture at p_k as the *replacement texture*. To generate an animation sequence of morphed images that depict the flow described by $\vec{V}(p)$ the texture advection algorithm is performed for all t_i , where $i = 0, \dots, n$. At each time step, if the integration ends prematurely, $t_k > t_0$ in Equation 1, then the particle has either reached the boundary or it has reached a critical point, where the velocity is zero. Again, note that for the case of steady flow, we can freely interchange t_k and t_i .

Suppose that the particle has reached the grid boundary, then two straightforward methods to determine the replacement texture for p_0 are as follows: (1) let the replacement texture be that at p_k where the integration ends; and (2) assume a periodic boundary condition such that the particle is wrapped around to the “opposite” side of the boundary. Both methods will produce undesirable artifacts. The first method would cause constant texture bands to flow into the image from the border, while the latter would cause incoherent textures to feed into the image from the border.

Ideally, we want to have a continuous texture to feed from the boundaries of the grid as time progresses. We found that the most effective method is to define a 3D procedural texture function $T(u, v, w)$ and compute the replacement texture on the fly whenever the particle integration ends. In Section 4, we will describe the texture function T in more detail. For now, we assume that the function T is given.

We determine the coordinates (u, v, w) of the 3D texture function based on the last position p_k where the integration ends and on the current t_k at the end of the integration. Let $u = \frac{p_{kx}}{M}$ and $v = \frac{p_{ky}}{N}$, where M and N are the image dimensions. If the particle integration ends prematurely, then we let $w = \frac{t_r}{t_n}$, where t_r is the remaining number of time steps that were not integrated i.e. $t_r = t_k - t_0$. Otherwise, we let $w = 0$ for all particle integrations that ends successfully. We can think of the texture function T being defined

in a unit cube and when the integration ends prematurely, we simply step into the cube along the w axis by a depth of t_r . Otherwise, we use $T(u, v, 0)$ as the replacement texture. Figure 2(a) illustrates this method.

3 FLOWS FROM REPELLING NODES

Here we are interested in generating flow textures from source critical points such that the physical flow behavior surrounding each critical point is clearly depicted. For this type of critical point, we want new textures to flow out of the critical point. We propose two approaches to generate such flow textures.

The first approach is based on the method described in the previous section for providing continuous texture feeding from the grid boundary. During particle advection, when the particle reaches the critical point, we determine the last position p_k (which most likely would be the critical point) and t_r , the remaining time steps to be integrated. Then, we compute the coordinates of (u, v, w) of the 3D procedural texture based on p_k and t_r . Though this approach would provide continuous texture being emitted from the source critical point, it will also produce concentric ring-like textures surrounding the critical point as time progresses. Figure 3 depicts the ring-like flow textures generated from the source critical points. Consider a group of particles that are at some equal distance from a source critical point. As these particles are traced backward, they would all reach the critical point at the same time. Hence, all would have the same replacement texture. Since this group of particles form a circle, we would get an expanding circular ring over time.

One may consider the ring-like texture patterns somewhat realistic because if one throws a pebble into a pool of water, one would see waves of rings from the position where the stone hit the water. However, the texture pattern does not really give the impression of the flow direction from the source critical points. Therefore, we propose another approach that achieves this effect. For the last position p_k , where the particle terminated, we step into the cube by a depth of t_r and then radially outward in the incoming direction of p_0 . Figure 2(b) illustrates this method. Let \vec{D} denote the incoming direction and $\vec{D} = \frac{p_0 - p_k}{\|p_0 - p_k\|}$, then we let $u = \frac{|p_k^x + R * D_x|}{M + R}$, $v = \frac{|p_k^y + R * D_y|}{N + R}$, and $w = \frac{t_r}{t_n}$, where R is a predefined constant. To keep the values of u and v between zero and one, we use the absolute value in the numerator because \vec{D} can be negative. With this approach, we were able to replace the ring-like effect on the animation with flow textures that radially emanate from the source critical points. Figure 4 depicts the results of this approach with $R = 180$.

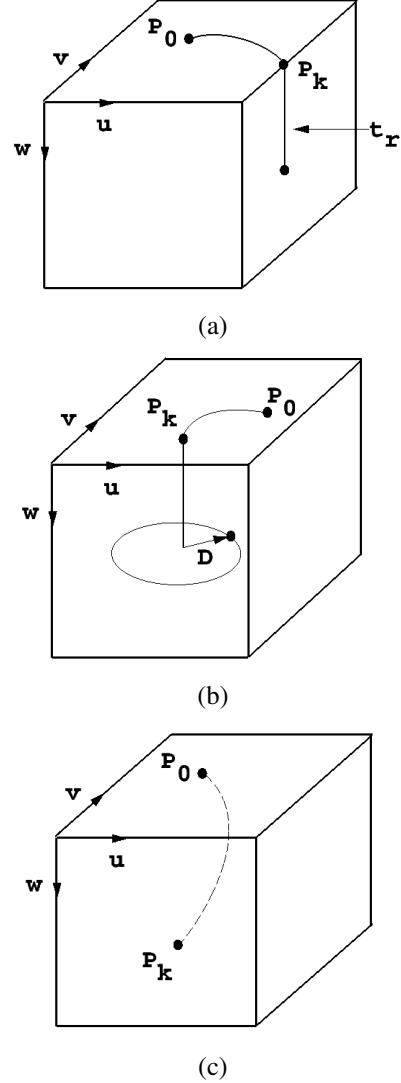
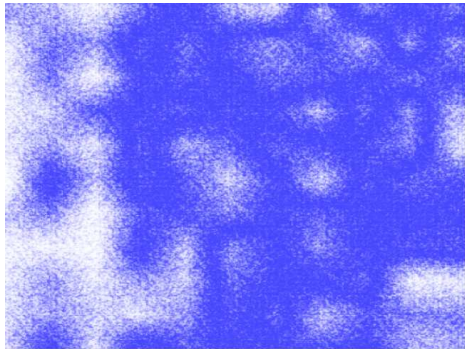
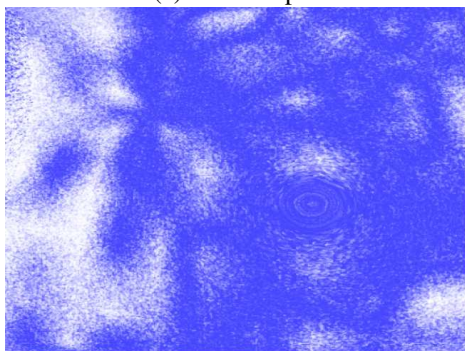


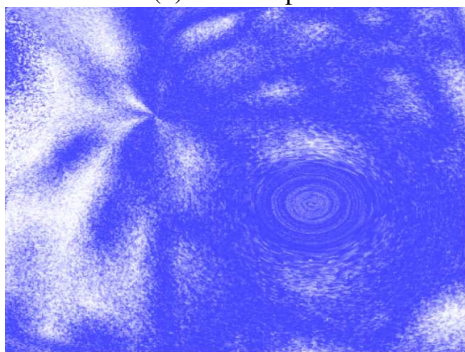
Figure 2. When particle reaches the boundary as shown in (a), the replacement texture is determined by stepping into the cube by a depth of t_r . If the particle reaches a source critical point as shown in (b), then the replacement texture is determined by stepping into the cube a depth of t_r and then radially outward in the \vec{D} direction. (c) depicts an alternative method where we advance into the texture volume as we advance our animation or integration, and does not require special handling of boundary inflow and source critical points.



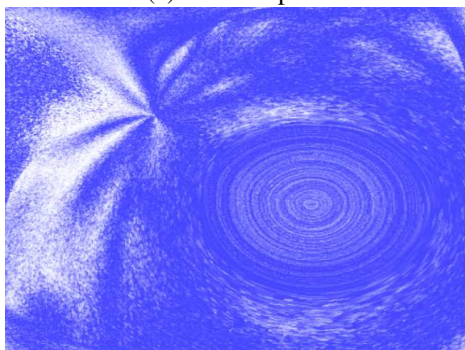
(a) Time Step 0



(b) Time Step 5

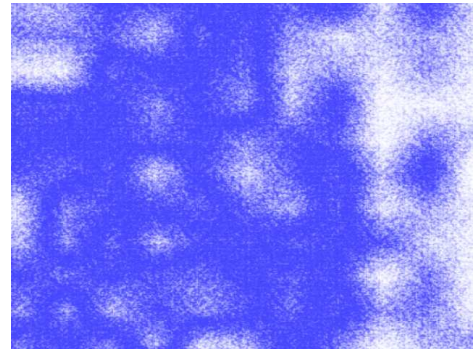


(c) Time Step 10

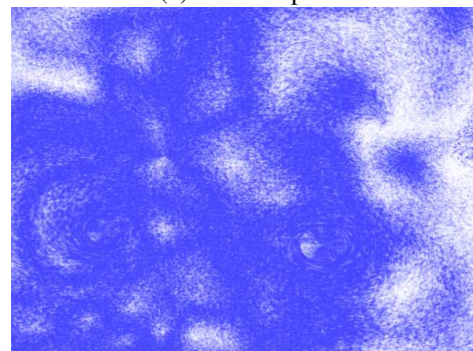


(d) Time Step 20

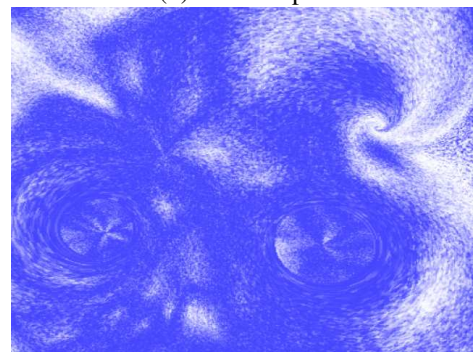
Figure 3. Flow data with one source critical point on the lower right and a sink critical point on the upper left. This time sequence depicts concentric ring-like texture patterns coming out of the source critical point.



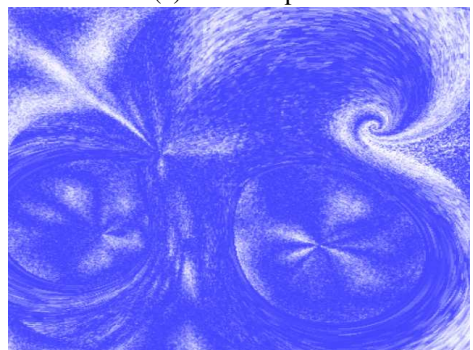
(a) Time Step 0



(b) Time Step 5



(c) Time Step 10



(d) Time Step 20

Figure 4. Flow data with two attracting nodes on the upper half and two repelling nodes on the lower half. This time sequence depicts flow textures emitting from the source critical points in the direction of the flow. The replacement texture is computed based on the method illustrated in Figure 2(b).

4 PROCEDURAL TEXTURE DESIGN

There are several issues and parameters to consider in designing textures for flow visualization. Poorly designed textures will create problems such as: (a) over-crowding of textural patterns particularly near attraction basins, and (b) lack of textural patterns in high shear regions. These can be addressed by first selecting a suitable spatial frequency of the textures to fit the complexity of the flow patterns. Over-crowding and absent of textures can be fixed by aging older textures and constantly introducing new textures. That is, we use 3D textures for 2D flow fields. Care must also be taken to balance the spatial frequency of the textures versus the temporal frequency so that the appearance of the advected textures does not alter the apparent speed of the flow.

We experimented with two types of solid textures. In both cases, the goal was to produce a cloud-like texture such that when it is advected and distorted, it would produce a wispy and fuzzy appearance of cloud being blown by wind. The first one is based on procedural textures as described in pages 155 and 156 of [1]. This technique uses a turbulence function to produce random looking clusters of various sizes and distributions. The texture value at any point can be quickly calculated on the fly and one does not need to store the 3D texture. Examples using this texture can be seen in Figures 3 and 4. However, one has little control over the distribution of the texture values, and in our application, may exacerbate the problem of over-crowding or under-population of advected textures over time.

The second technique creates 3D textures by populating a volume with spheres. The parameters include average size of the spheres, average distance between spheres, and how the texture varies within a sphere. We use a poisson disk distribution to maintain a minimum average distance among the spheres. For each sphere, we vary the texture value from 1.0 at the center to 0.5 at the surface. This helps produce softer fuzzy spheres.

In generating the 3D textures, we note that one must take into account both spatial and temporal scales of the flow data. That is, the patterns in the texture must not be too large so that it misses the flow pattern. For example, if the texture consists of one very large sphere and the flow consist of numerous critical points, then the advected texture will not be able to effectively reveal the nature of the flow. Likewise, the patterns in the texture must not be too small so that it quickly disappears and fades without showing the flow pattern. With regards to temporal scales, the patterns should not vary too slowly or they linger around too long in the animation and create over-crowding of textures; neither should they vary too fast or they become distractions that make the animation blink with very short lived flashes. Figure 5 depicts a time-sequence from the animation of three

rotating vortices. In the animation, we advance into the texture volume as illustrated in Figure 2(c). In Figure 5(a), the undistorted spheres in the texture are shown. In (b), we start to see the distortion of the spheres caused by two rotating vortices near the upper right and lower left of the image. In (c), the texture distortions caused by central vortex is also apparent. Finally, in (d), the effects of the three vortices on the texture are shown. In addition, there are two saddle structures located near the upper left and lower right of the image.

Using the sphere texture volume, we found the animation generated by advancing into the texture volume at each time step (Figure 2(c)) is more sensitive to the size and number of spheres in the volume. We had to experiment with several sphere sizes to obtain an animation that effectively revealed the flow. On the other hand, the frequency scales included in the solid textures based on turbulence functions seemed to be more robust.

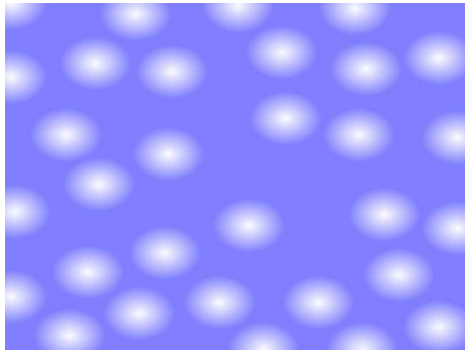
Currently, we generate our textures based on trial and error, and do not yet have a systematic nor heuristic method for automatically determining the appropriate spatial and temporal scales given a flow field. The scales obviously will depend on the magnitude and complexity of the flow, and we plan to pursue an investigation on how to generate a flow dependent texture for animation.

5 RESULTS

We now describe some results of applying the proposed techniques on two CFD data sets. The first data set is the first time step of a cross section of flow around an oscillating airfoil. A good background description of the data can be found in [4, 6]. The second data set is surface flow information on a hemisphere cylinder. Alternative ways of visualizing this data set can be found in [9, 2]

Figure 6 shows four frames of our flow animation using texture advection. In (a), we see the geometry of the wing profile, how the procedural texture looks like before it is advected, and also the discontinuity at the trailing tip of the wing due to the wrap around of the curvilinear grid and non-periodic textures. In (b), we start to see the formation of a primary clockwise vortex above the airfoil, and a secondary counter-clockwise vortex near the trailing tip of the wing. In (c), the primary vortex becomes more prominent. At the same time, we start to see some compression of textures forming at the leading edge of the wing. Finally, in (d), the primary vortex has separated from the airfoil and the compressed textures at the leading edge have become even more prominent. This sequence is best seen in an animation which is available for download from www.cse.ucsc.edu/research/avis/textflow2d.html.

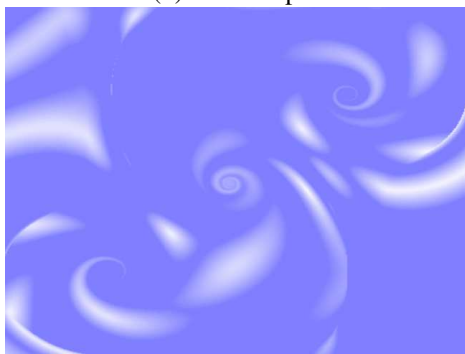
Figure 7 shows three views of three different time steps in our texture advection animation of the surface flow on



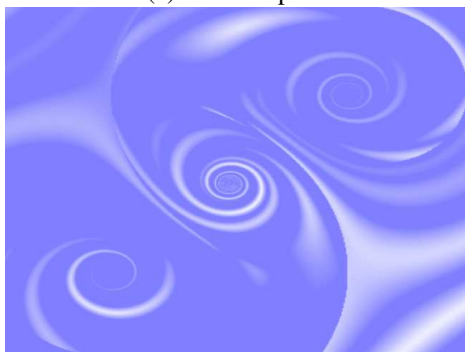
(a) Time Step 0



(b) Time Step 5

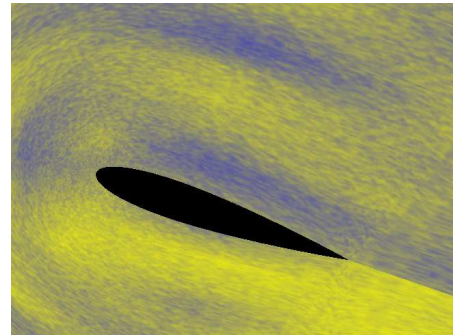


(c) Time Step 16

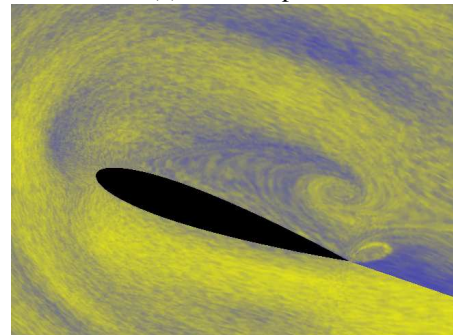


(d) Time Step 45

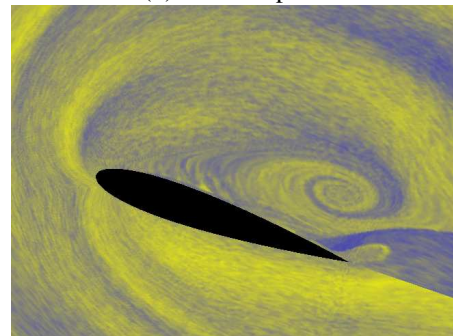
Figure 5. Advection of 3D sphere textures using flow data with three rotating vortices located diagonally across the image.



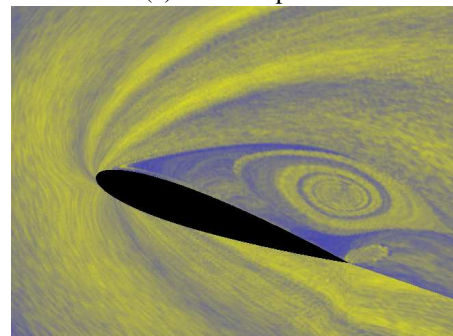
(a) Time Step 0



(b) Time Step 10



(c) Time Step 25



(d) Time Step 60

Figure 6. Data is from a cross sectional slice of a CFD computation for flow around an air-foil. More description of this data can be found in [4, 6]. One can see the formation of vortices above and behind the airfoil, as well as compression of textures off the leading edge of the foil.

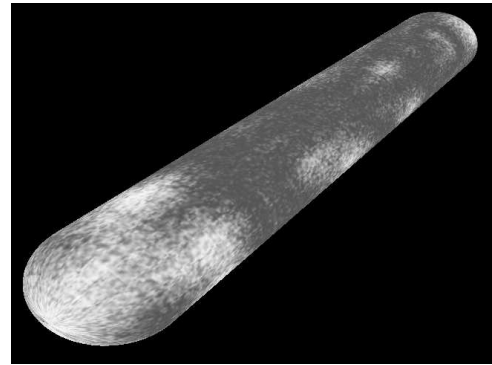
the hemisphere cylinder. In (a), we see the initial mapping of the un-advected textures in the physical domain of the cylinder. In (b), we see a source critical point near the top of the surface. Also, the sharp discontinuous curve running through most of the length of the cylinder correspond to surface topology lines and indicates regions of flow with very different directions. In (c), the animation has progressed further along and shows the persistence of certain flow patterns even though the textures are dynamically being advected. Upon closer examination, one can see two separation lines along the cylinder body. Note that there is also another critical point between the two separation lines. Furthermore, there is a source critical point near the nose of the cylinder. Again, these are best observed in an animation which is available from the same web site mentioned above.

When we first showed these texture animation of real flow data to our CFD scientists, they were very surprised with seeing the animation. They have never seen this type of synthetic texture flow on their flow data. Previously, they had mostly relied on the traditional streamline particle traces to observe their flow trajectory. Almost all of them have very good comments about the animation and would like to see our technique available to them soon.

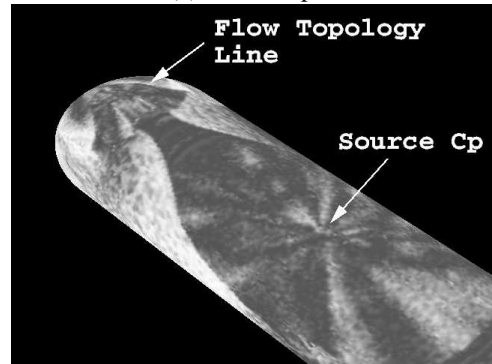
6 CONCLUSION

We have described techniques and textures for animating 2D steady flow fields using texture advection. The method is straight forward and relatively easy to implement. The modifications introduced in this paper also addressed all the difficulties identified in earlier texture advection work – over-crowding and/or sparsity of textures over time, and handling of boundary inflow and repelling nodes. The paper first demonstrated how the proposed methods and textures handled problematic cases using synthetic flow data. We then applied it to two CFD data sets. Looking at still images in Figures 6 and 7, one can already gain the benefit of this technique in the form of automatic extraction of flow features such as location and type of critical points, as well as flow topology curves. These features and the flow pattern in general are further accentuated when one views the animations.

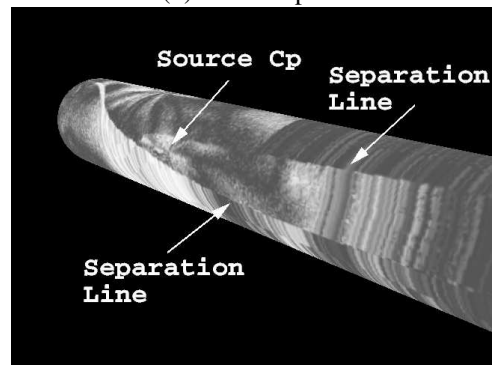
There are a number of places for improvement and extension: (1) we think texture advection will be most beneficial for unsteady flow data – this is what we observe in nature as clouds form and dissipate, or as dyes swirl in a mixture; (2) the selection of the spatial and temporal scales of our textures is currently by trial and error, we hope to come up with heuristics or recommendations for their selection based on information about flow magnitudes; (3) the patterns coming out of source critical points may be either concentric ring-like patterns as one might observe when a pebble is dropped in a pool, or it may be appear like pat-



(a) Time Step 0



(b) Time Step 17



(c) Time Step 34

Figure 7. Surface flow data on the hemisphere cylinder. Critical points and topology lines become obvious as the animation proceeds.

terns shooting out of the source. When one is the mind-set of the latter case, we still have a distracting ring pattern. We are exploring different ways of eliminating this extra ring; and (4) extension of this technique to steady and time varying 3D flow fields. There is also excellent work in texture advection of flow data carried out by Bruno Jobard that uses simple white noise textures.¹

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References

- [1] David S. Ebert, F. Kenton Musgrave, Darwyn Peachey, Ken Perlin, and Steven Worley. *Texturing and Modeling, A Procedural Approach*. AP Professional, 1994.
- [2] J. L. Helman and Lambertus Hesselink. Surface representations of two and three-dimensional fluid flow topology. In *Proceedings: Visualization '90*, pages 6–13. IEEE Computer Society, 1990.
- [3] David Kao and Alex Pang. On animating 2D velocity fields. In *SPIE Visual Data Exploration and Analysis VIII*, volume 4302, pages 41–48, 2001.
- [4] S. Ko and W. McCrosky. Computations of unsteady separating flows over an oscillating airfoil. *AIAA, 33rd Aerospace Sciences Meeting*, January 1995. Paper number: 95-0312.
- [5] N. Max, R. Crawfis, and D. Williams. Visualizing wind velocities by advecting cloud textures. In *Proceedings: Visualization '92*, pages 179–184. IEEE Computer Society, 1992.
- [6] Han-Wei Shen and David L. Kao. A new line integral convolution algorithm for visualizing time-varying flow fields. *IEEE Transactions on Visualization and Computer Graphics*, 4(2):98–108, April-June 1998.
- [7] Jos Stam. Stable fluids. In *Computer Graphics*, pages 121–128, Los Angeles, Ca., August 1999. ACM Siggraph Conference Proceedings.
- [8] Patrick Witting. Computational fluid dynamics in a traditional animation environment. In *Computer Graphics*, pages 129–136, Los Angeles, Ca., August 1999. ACM Siggraph Conference Proceedings.
- [9] S. X. Ying and J. L. Steger L. B. Schiff. A numerical study of three-dimensional separated flow past a hemisphere cylinder. In *Proceedings of AIAA 19th Fluid Dynamics, Plasma Dynamics and Lasers Conference*. American Institute of Aeronautics and Astronautics, June 1987. Paper number: 87-1207.

¹Based on personal communication with Bruno Jobard. See www.csit.fsu.edu/~jobard/SwissMeteo/meteo.html.